

ENERGY POSITION OF MAXIMA IN RESPONSE FUNCTIONS OF ^2H NUCLEUS WITH $q = 1 - 1,5 \text{ fm}^{-1}$

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The problem of measuring a quasielastic electron scattering peak energy on nuclei is discussed.

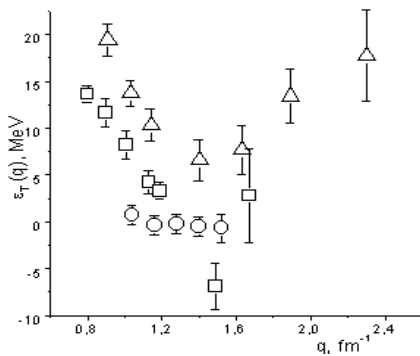
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The energy of a peak maximum in quasielastic electron scattering (QES) on nuclear nucleons is designated as ε and is measured from the energy of elastic electron scattering on the nucleon with the same transfer 3-momenta q , that corresponds to a maximum of this peak. The theoretical works [1,2] and some others show that the calculation of $\varepsilon(q)$ depends on the choice of a nucleon potential type. Some of calculated $\varepsilon(q)$ functions have a minimum for $q < 1.5 \text{ fm}^{-1}$, where some minimum in data of ε measurements on the ^{12}C nucleus [3] is observed too. But the result interpretation of such data in the aspect of [1,2] is faced with two difficulties:

I. Originally the experimental ε were defined from dependencies of twice differential cross sections $d^2\sigma$ on the transfer energy ω (values ε obtained by this way we designate as ε_σ). This cross section is a sum of longitudinal R_L and transverse R_T response functions:

$$d^2\sigma \cong \sigma_M(G(q))^2 \{R_L(q, \omega) + [0.5 + tg(\theta/2)]R_T(q, \omega)\},$$

where σ_M is the Mott cross section, $G(q)$ is the proton electric form factor, θ is the angle of electron scattering. It was showed (e.g., see [2]) that ε_L , corresponding to the R_L -function and ε_T - R_T -function, are unequal. So, according to the equation for $d^2\sigma$, the value ε_σ is in the interval $\varepsilon_L - \varepsilon_T$ and the result of the ε_σ measurement depends on the choice of the angle θ .



II. With the reducing q the QES peak is displaced on ω into the area of excitation of nuclear bound states (NBS), such as giant resonances, and then the experimental ε correspond to the sum peak maximum of nuclear electrodisintegration but it is not a peak of QES described by theoretical calculations.

The experiment has shown that really $\varepsilon_L \neq \varepsilon_T$ [4], and the dependence of ε_T on q at $q \leq 1.3 \text{ fm}^{-1}$ is completely described by the calculation in which the function $\varepsilon(q)$ is specified by NBS excitation [5,6]. Analysis of these results in view of the giant resonance systematic leads to

the conclusion, that for small transfer momentum the measurement of ε_L and ε_T , which characterize QES, is rather difficult. The ^2H nucleus is interesting because it has no excitation states and thus there is a possibility to study ε of QES at small q experimentally.

In the present paper we use the results of measurement of the R_L - and R_T -functions of the ^2H nucleus at the linac LUE-300 KIPT and obtain the ε values for $q = 1 - 1.5 \text{ fm}^{-1}$. We have determined for the ^2H nucleus $\varepsilon_L - \varepsilon_T \approx 1 \text{ MeV}$. The obtained ε_T values are shown in the figure by circles. Also, for comparison, on the figure the ε_T -data from [5] and [6] for ^6Li (triangles) and ^4He (squares) respectively are quoted. One can see that at $q < 1.5 \text{ fm}^{-1}$ for ε of the ^2H nucleus there is no the dependence on q , that is character for ε of other nuclei. This difference in the behavior of dependences of ε on q is a confirmation of NBS role in shaping the maximum peak of nuclear electrodisintegration.

So obtained experimental ε of the ^2H nucleus allow extending an area of the calculations checking of a peak QES energy position at q less than 1.5 fm^{-1} .

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