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THE PROBLEM OF CONTROL OF MEMBRANE VIBRATIONS WITH NON-SEPARATED MULTIPOINT CONDITIONS AT INTERMEDIATE MOMENTS OF TIME

Introduction. Many control processes from various fields of science and technology lead to the necessity to study multipoint boundary value problems of control, in which, along with classical boundary conditions, non-separated multi-point intermediate conditions are also given. A characteristic feature of multipoint boundary value problems of control is the presence of non-separated conditions at several intermediate points of the study interval. Such control problems have important applied and theoretical value, a necessity naturally arises for their investigation in various settings. In this article, the problem of control of vibrations of a rectangular membrane with given initial, final conditions and non-separated values of the deflection function and velocities at intermediate moments of time is considered.

The purpose of the paper is to develop a constructive approach to construct a function of control action to control the vibrations of a rectangular membrane with given initial, final conditions and non-separated (non-local) values of the deflection and velocities of membrane points at intermediate moments of time.

Results. By the method of separation of variables, the problem is reduced to the problem of control of ordinary differential equations with given initial, final, and non-separated multipoint intermediate conditions. Using the methods of the theory of control of finite-dimensional systems with multipoint intermediate conditions, a control action to control vibrations of a rectangular membrane is constructed.

Conclusion. The problem of control of the vibrations of a rectangular membrane with given non-separated values of the deflection function and velocities at intermediate moments of time is solved by using the methods of the theory of control of finite-dimensional systems with multipoint intermediate conditions.

Keywords: control of vibrations, membrane vibration, intermediate values, non-separated multipoint conditions.

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INTRODUCTION

One of the most common processes in nature and technology are oscillatory processes, which are modeled by the wave equations [1–4]. At the same time, in practice, control problems often arise when it is necessary to generate the desired vibration form that satisfies intermediate conditions. Many control processes lead to the necessity to study multipoint boundary value problems, in which, along with the classical boundary (initial and final) conditions, non-separated (non-local) multipoint intermediate conditions are given [5–15]. Non-separated multipoint boundary value problems, on the one hand, arise as mathematical models of real processes, and on the other hand, for many equations correct formulation of local boundary value problems is impossible. The nonseparability of multipoint conditions may be, in particular, due to the impossibility in practice to measure the parameters of the state of an object instantaneously or at its individual points.

Numerous examples of technological processes leading to the problems of control of systems with distributed parameters were considered in [1–3] and various methods of solutions were proposed. The control problems of oscillatory processes, with both external and boundary control actions under various types of boundary conditions are considered in [4, 7–14]. In [7–12] problems of control of vibrations of a string and membrane with given intermediate (local) states with the help of external forces acting on the systems were considered. The problems of control of distributed systems with given non-separated multipoint (non-local) conditions at intermediate moments of time are less explored to date.

The purpose of the paper is to develop a constructive approach to construct a function of control action for controlling vibrations of a rectangular membrane with given initial, final conditions and non-separated (non-local) deflection values and velocities of membrane points at intermediate times.

THE FORMULATION OF THE PROBLEM

Consider a homogeneous, elastic, rectangular membrane, the edges of which are fixed. Let distributed forces act on the membrane with a density $u(x, y, t)$ perpendicular to the surface of the membrane, under the action of which the membrane will vibrate. We restrict ourselves to the consideration of small vibrations of the membrane.

The state of the membrane is described by the function $Q(x, y, t)$, $0 \leq x \leq b$, $0 \leq y \leq c$, $0 < t < T$, which for $0 < x < b$, $0 < y < c$ и $0 < t < T$ is characterized by the following equation:

$$\frac{\partial^2 Q}{\partial t^2} = a^2 \left(\frac{\partial^2 Q}{\partial x^2} + \frac{\partial^2 Q}{\partial y^2} \right) + u(x, y, t) \quad (1)$$

with homogeneous boundary conditions

$$Q(0, y, t) = 0, Q(b, y, t) = 0, Q(x, 0, t) = 0, Q(x, c, t) = 0, 0 \leq t \leq T \quad (2)$$

and satisfying initial and final conditions

$$Q(x, y, 0) = \varphi_0(x, y), \quad \left. \frac{\partial Q}{\partial t} \right|_{t=0} = \psi_0(x, y), \quad 0 \leq x \leq b, \quad 0 \leq y \leq c, \quad (3)$$

$$Q(x, y, T) = \varphi_T(x, y) = \varphi_{m+1}(x, y), \quad \left. \frac{\partial Q}{\partial t} \right|_{t=T} = \psi_T(x, y) = \psi_{m+1}(x, y), \quad (4)$$

$$0 \leq x \leq b, \quad 0 \leq y \leq c.$$

On the left hand side of equation (1) the function $u(x, y, t)$ - density of force, which is the control action, $a^2 = \frac{T_0}{\rho}$, where T_0 - tension, and ρ - membrane density. It is assumed that the function $u(x, y, t) \in L_2(\Omega)$, where $\Omega = \{(x, y, t) : x \in [0, b], y \in [0, c], t \in [0, T]\}$.

Let at some intermediate moments of time $0 = t_0 < t_1 < \dots < t_m < t_{m+1} = T$ on the values of the membrane deflection function non-separated (non-local) conditions are given in the following form:

$$\sum_{k=1}^m f_k Q(x, y, t_k) = \alpha(x, y), \quad (5)$$

$$\sum_{k=1}^m e_k \left. \frac{\partial Q(x, y, t)}{\partial t} \right|_{t=t_k} = \beta(x, y), \quad (6)$$

where f_k and e_k — given values ($k = 1, \dots, m$), $\alpha(x, y)$ and $\beta(x, y)$ — some known functions.

Necessary condition for existence of the classical solution of the formulated problem (1)–(6) is a condition of compatibility of boundary, initial, intermediate and final conditions. Therefore it is supposed that $\varphi_0(x, y)$, $\psi_0(x, y)$, $\varphi_T(x, y)$, $\psi_T(x, y)$, $\alpha(x, y)$ and $\beta(x, y)$ are given smooth functions satisfying the compatibility conditions.

In general, it is possible that at some moments of time t_k ($k = 1, \dots, m$) in the conditions (5), (6) either the values of the deflection function or values of derivatives of that function are present, i.e. it is not mandatory that at each moment of time t_k ($k = 1, \dots, m$) in the conditions (5), (6) functions $Q(x, y, t_k)$ and $\left. \frac{\partial Q(x, y, t)}{\partial t} \right|_{t=t_k}$ are present at the same time. In those case we will consider that

the corresponding coefficients f_k or e_k are equal to zero. In particular, assuming that $f_2 = e_1 = 0$ and $f_1 = e_2 = 1$, conditions (5) and (6) take the following form:

$$Q(x, y, t_1) = \alpha(x, y), \quad \left. \frac{\partial Q(x, y, t)}{\partial t} \right|_{t=t_2} = \beta(x, y).$$

The control problem of membrane vibrations with given non-separated values of the deflection function and velocities at intermediate moments of time t_k ($k=1, \dots, m$) can be formulated as follows: among the possible controls $u(x, y, t)$, $0 \leq x \leq b, 0 \leq y \leq c, 0 \leq t \leq T$ it is required to find a control that transfers membrane vibrations (1) with boundary conditions (2) from a given initial state (3) to a given final state (4), ensuring the satisfaction of non-separated multipoint intermediate conditions (5), (6).

It is assumed that system (1) with constraints (2) - (6) over the time interval $[0, T]$ is completely controllable [5, 16].

THE SOLUTION OF THE PROBLEM

To build a solution to the formulated problem, we are looking for a solution to the equation (1) with boundary conditions (2) in the following form:

$$Q(x, y, t) = \sum_{k,n=1}^{\infty} Q_{kn}(t) \sin \frac{k\pi}{b} x \sin \frac{n\pi}{c} y. \quad (7)$$

Apparently, to determine $Q(x, y, t)$ it suffices to determine $Q_{kn}(t)$, where $k, n = 1, 2, \dots$. We represent the function $u(x, y, t)$ in the form of Fourier series

$$u(x, y, t) = \sum_{k,n=1}^{\infty} u_{kn}(t) \sin \frac{k\pi}{b} x \sin \frac{n\pi}{c} y. \quad (8)$$

Let us substitute decomposition (7), (8) into the relation (1). By virtue of the orthogonality of the system of eigenfunctions $\left\{ \sin \frac{k\pi}{b} x \sin \frac{n\pi}{c} y \right\}$, where $k, n = 1, 2, \dots$, it follows that coefficients of Fourier $Q_{kn}(t)$ satisfy the countable system of ordinary differential equations

$$\ddot{Q}_{kn}(t) + \lambda_{kn}^2 Q_{kn}(t) = u_{kn}(t), \quad \lambda_{kn}^2 = a^2 \left[\left(\frac{k\pi}{b} \right)^2 + \left(\frac{n\pi}{c} \right)^2 \right], \quad (9)$$

$$k, n = 1, 2, \dots$$

and the following initial, non-separated multipoint intermediate and final conditions:

$$Q_{kn}(0) = \varphi_{kn}^{(0)}, \quad \dot{Q}_{kn}(0) = \psi_{kn}^{(0)}, \quad (10)$$

$$\sum_{j=1}^m f_j Q_{kn}(t_j) = \alpha_{kn}, \quad \sum_{j=1}^m e_j \dot{Q}_{kn}(t_j) = \beta_{kn}, \quad (11)$$

$$Q_{kn}(T) = \varphi_{kn}^{(T)} = \varphi_{kn}^{(m+1)}, \quad \dot{Q}_{kn}(T) = \psi_{kn}^{(T)} = \psi_{kn}^{(m+1)}, \quad (12)$$

where $Q_{kn}(t)$, $\varphi_{kn}^{(0)}$, $\psi_{kn}^{(0)}$, $\varphi_{kn}^{(m+1)}$, $\psi_{kn}^{(m+1)}$, $u_{kn}(t)$, α_{kn} and β_{kn} denote Fourier coefficients, corresponding to the functions $Q(x, y, t)$, $\varphi_0(x, y)$, $\psi_0(x, y)$, $\varphi_{m+1}(x, y)$, $\psi_{m+1}(x, y)$, $u(x, y, t)$, $\alpha(x, y)$ and $\beta(x, y)$.

The general solution of equation (9) with initial conditions (10) and its time derivative have the following form:

$$Q_{kn}(t) = \varphi_{kn}^{(0)} \cos \lambda_{kn} t + \frac{1}{\lambda_{kn}} \psi_{kn}^{(0)} \sin \lambda_{kn} t + \frac{1}{\lambda_{kn}} \int_0^t u_{kn}(\tau) \sin \lambda_{kn}(t - \tau) d\tau, \quad (13)$$

$$\dot{Q}_{kn}(t) = -\lambda_{kn} \varphi_{kn}^{(0)} \sin \lambda_{kn} t + \psi_{kn}^{(0)} \cos \lambda_{kn} t + \int_0^t u_{kn}(\tau) \cos \lambda_{kn}(t - \tau) d\tau.$$

Now, taking into account the intermediate non-separated (11) and final (12) conditions using the approaches given in [6, 7], from equation (13), we obtain that the functions $u_{kn}(\tau)$ for each k and n must satisfy the following system of equalities:

$$\begin{aligned} \int_0^T u_{kn}(\tau) \sin \lambda_{kn}(T - \tau) d\tau &= C_{1kn}(T), \\ \int_0^T u_{kn}(\tau) \cos \lambda_{kn}(T - \tau) d\tau &= C_{2kn}(T), \\ \sum_{j=1}^m f_j \int_0^{t_j} u_{kn}(\tau) \sin \lambda_{kn}(t_j - \tau) d\tau &= C_{1kn}^{(m)}(t_1, \dots, t_m), \\ \sum_{j=1}^m e_j \int_0^{t_j} u_{kn}(\tau) \cos \lambda_{kn}(t_j - \tau) d\tau &= C_{2kn}^{(m)}(t_1, \dots, t_m), \end{aligned} \quad (14)$$

where

$$\begin{aligned} C_{1kn}(T) &= \lambda_{kn} \varphi_{kn}^{(m+1)} - \lambda_{kn} \varphi_{kn}^{(0)} \cos \lambda_{kn} T - \psi_{kn}^{(0)} \sin \lambda_{kn} T, \\ C_{2kn}(T) &= \psi_{kn}^{(m+1)} + \lambda_{kn} \varphi_{kn}^{(0)} \sin \lambda_{kn} T - \psi_{kn}^{(0)} \cos \lambda_{kn} T, \end{aligned} \quad (15)$$

$$C_{1kn}^{(m)}(t_1, \dots, t_m) = \lambda_{kn} \left[\alpha_{kn} - \sum_{j=1}^m f_j \left(\varphi_{kn}^{(0)} \cos \lambda_{kn} t_j + \frac{1}{\lambda_{kn}} \psi_{kn}^{(0)} \sin \lambda_{kn} t_j \right) \right],$$

$$C_{2kn}^{(m)}(t_1, \dots, t_m) = \beta_{kn} - \sum_{j=1}^m e_j \left(-\lambda_{kn} \varphi_{kn}^{(0)} \sin \lambda_{kn} t_j + \psi_{kn}^{(0)} \cos \lambda_{kn} t_j \right).$$

We introduce the following functions

$$h_{1kn}(\tau) = \sin \lambda_{kn}(T - \tau), \quad h_{2kn}(\tau) = \cos \lambda_{kn}(T - \tau), \quad 0 \leq \tau \leq T$$

$$h_{1kn}^{(m)}(\tau) = \sum_{j=1}^m f_j h_{1kn}^{(j)}(\tau)$$

$$h_{1kn}^{(j)}(\tau) = \begin{cases} \sin \lambda_{kn}(t_j - \tau) & \text{when } 0 \leq \tau \leq t_j, \\ 0 & \text{when } t_j < \tau \leq t_{m+1} = T, \end{cases} \quad (16)$$

$$h_{2kn}^{(m)}(\tau) = \sum_{j=1}^m e_j h_{2kn}^{(j)}(\tau),$$

$$h_{2kn}^{(j)}(\tau) = \begin{cases} \cos \lambda_{kn}(t_j - \tau) & \text{when } 0 \leq \tau \leq t_j, \\ 0 & \text{when } t_j < \tau \leq t_{m+1} = T. \end{cases}$$

Then the integral relations (14) with the help of the function (16) are written as follows:

$$\int_0^T u_{kn}(\tau) h_{1kn}(\tau) d\tau = C_{1kn}(T), \quad \int_0^T u_{kn}(\tau) h_{2kn}(\tau) d\tau = C_{2kn}(T),$$

$$\int_0^T u_{kn}(\tau) h_{1kn}^{(m)}(\tau) d\tau = C_{1kn}^{(m)}(t_1, \dots, t_m), \quad \int_0^T u_{kn}(\tau) h_{2kn}^{(m)}(\tau) d\tau = C_{2kn}^{(m)}(t_1, \dots, t_m), \quad (17)$$

$$k, n = 1, 2, \dots$$

Thus, the required functions $u_{kn}(\tau)$, $\tau \in [0, T]$, for each k and n must satisfy the integral relations (17).

Using the following notation

$$H_{kn}(\tau) = \begin{pmatrix} h_{1kn}(\tau) \\ h_{2kn}(\tau) \\ h_{1kn}^{(m)}(\tau) \\ h_{2kn}^{(m)}(\tau) \end{pmatrix}, \quad \eta_{kn} = \begin{pmatrix} C_{1kn}(T) \\ C_{2kn}(T) \\ C_{1kn}^{(m)}(t_1, \dots, t_m) \\ C_{2kn}^{(m)}(t_1, \dots, t_m) \end{pmatrix} \quad (18)$$

we re-write integral relations (17) in the form:

$$\int_0^T H_{kn}(t) u_{kn}(t) dt = \eta_{kn}. \quad (19)$$

From relation (19) (or (17)) it follows that for each harmonic, the motion described by equation (9) with conditions (10) - (12) is completely controllable if and only if for any given vector η_{kn} (18) it is possible to find a control $u_{kn}(t)$, $t \in [0, T]$, satisfying condition (19) (or (17)).

We introduce notation

$$S_{kn} = \int_0^T H_{kn}(t)(H_{kn}(t))^T dt = \begin{pmatrix} S_{11}^{(kn)} & S_{12}^{(kn)} & S_{13}^{(kn)} & S_{14}^{(kn)} \\ S_{21}^{(kn)} & S_{22}^{(kn)} & S_{23}^{(kn)} & S_{24}^{(kn)} \\ S_{31}^{(kn)} & S_{32}^{(kn)} & S_{33}^{(kn)} & S_{34}^{(kn)} \\ S_{41}^{(kn)} & S_{42}^{(kn)} & S_{43}^{(kn)} & S_{44}^{(kn)} \end{pmatrix}, \quad (20)$$

where $H_{kn}(t)(H_{kn}(t))^T$ – outer product of vectors. Henceforth the letter « T » in the superscript denotes transpose operation.

Let's assume that $\det S_n \neq 0$. Then, following [5, 17], for each $k, n = 1, 2, \dots$ the function $u_{kn}(t)$, $t \in [0, T]$, satisfying the integral relation (19), can be written as

$$u_{kn}(t) = (H_{kn}(t))^T S_{kn}^{-1} \eta_{kn} + v_{kn}(t), \quad (21)$$

where $v_{kn}(t)$ – some vector function such that

$$\int_0^T H_{kn}(t)v_{kn}(t)dt = 0. \quad (22)$$

The elements of the matrix S_n , according to (20) and the notation (16), (18), have the following forms:

$$\begin{aligned} s_{11}^{(kn)} &= \int_0^T (h_{1kn}(\tau))^2 d\tau = \int_0^T (\sin \lambda_{kn}(T-\tau))^2 d\tau, \\ s_{12}^{(kn)} &= s_{21}^{(kn)} = \int_0^T h_{1kn}(\tau)h_{2kn}(\tau)d\tau = \int_0^T \sin \lambda_{kn}(T-\tau) \cos \lambda_{kn}(T-\tau)d\tau, \\ s_{13}^{(kn)} &= s_{31}^{(kn)} = \int_0^T h_{1kn}(\tau)h_{1kn}^{(m)}(\tau)d\tau = \int_0^T \sin \lambda_{kn}(T-\tau) \left(\sum_{j=1}^m f_j h_{1kn}^{(j)}(\tau) \right) d\tau, \\ s_{14}^{(kn)} &= s_{41}^{(kn)} = \int_0^T h_{1kn}(\tau)h_{2kn}^{(m)}(\tau)d\tau = \int_0^T \sin \lambda_{kn}(T-\tau) \left(\sum_{j=1}^m e_j h_{2kn}^{(j)}(\tau) \right) d\tau = \\ &= \sum_{j=1}^m e_j \int_0^T \sin \lambda_{kn}(T-\tau) h_{2kn}^{(j)}(\tau) d\tau = \sum_{j=1}^m e_j \int_0^{t_j} \sin \lambda_{kn}(T-\tau) \cos \lambda_{kn}(t_j-\tau) d\tau, \\ s_{23}^{(kn)} &= s_{32}^{(kn)} = \int_0^T h_{2kn}(\tau)h_{1kn}^{(m)}(\tau)d\tau = \int_0^T \cos \lambda_{kn}(T-\tau) \left(\sum_{j=1}^m f_j h_{1kn}^{(j)}(\tau) \right) d\tau, \end{aligned} \quad (23)$$

$$\begin{aligned}
 s_{22}^{(kn)} &= \int_0^T (h_{2kn}(\tau))^2 d\tau = \int_0^T (\cos \lambda_{kn}(T-\tau))^2 d\tau, \\
 s_{24}^{(kn)} &= s_{42}^{(kn)} = \int_0^T h_{2kn}(\tau) h_{2kn}^{(m)}(\tau) d\tau = \int_0^T \cos \lambda_{kn}(T-\tau) \left(\sum_{j=1}^m e_j h_{2kn}^{(j)}(\tau) \right) d\tau = \\
 &= \sum_{j=1}^m e_j \int_0^T \cos \lambda_{kn}(T-\tau) h_{2kn}^{(j)}(\tau) d\tau = \sum_{j=1}^m e_j \int_0^{t_j} \cos \lambda_{kn}(T-\tau) \cos \lambda_{kn}(t_j-\tau) d\tau, \\
 s_{33}^{(kn)} &= \int_0^T (h_{1kn}^{(m)}(\tau))^2 d\tau = \int_0^T \left(\sum_{j=1}^m f_j h_{1kn}^{(j)}(\tau) \right)^2 d\tau, \\
 s_{34}^{(kn)} &= s_{43}^{(kn)} = \int_0^T h_{1kn}^{(m)}(\tau) h_{2kn}^{(m)}(\tau) d\tau = \int_0^T \left(\sum_{j=1}^m f_j h_{1kn}^{(j)}(\tau) \right) \left(\sum_{j=1}^m e_j h_{2kn}^{(j)}(\tau) \right) d\tau, \\
 s_{44}^{(kn)} &= \int_0^T (h_{2kn}^{(m)}(\tau))^2 d\tau = \int_0^T \left(\sum_{j=1}^m e_j h_{2kn}^{(j)}(\tau) \right)^2 d\tau.
 \end{aligned}$$

Note that, according to the notation (16), we will have

$$h_{1kn}^{(m)}(t) = \begin{cases} \sum_{j=1}^m f_j \sin \lambda_{kn}(t_j - t), & 0 \leq t \leq t_1 \\ \sum_{j=2}^m f_j \sin \lambda_{kn}(t_j - t), & t_1 < t \leq t_2 \\ \dots & \\ \sum_{j=m-1}^m f_j \sin \lambda_{kn}(t_j - t), & t_{m-2} < t \leq t_{m-1} \\ f_m \sin \lambda_{kn}(t_m - t), & t_{m-1} < t \leq t_m \\ 0, & t_m < t \leq t_{m+1} = T \end{cases},$$

$$h_{2kn}^{(m)}(t) = \begin{cases} \sum_{j=1}^m e_j \cos \lambda_{kn}(t_j - t), & 0 \leq t \leq t_1 \\ \sum_{j=2}^m e_j \cos \lambda_{kn}(t_j - t), & t_1 < t \leq t_2 \\ \dots & \\ \sum_{j=m-1}^m e_j \cos \lambda_{kn}(t_j - t), & t_{m-2} < t \leq t_{m-1} \\ e_m \cos \lambda_{kn}(t_m - t), & t_{m-1} < t \leq t_m \\ 0, & t_m < t \leq t_{m+1} = T \end{cases}.$$

Therefore, taking into account the notation (16) and (18), the control action $u_{kn}(t)$, $t \in [0, T]$, according to (21), is represented in the following form:

$$u_{kn}(t) = \begin{cases} \left(\sin \lambda_{kn}(T-t) \cos \lambda_{kn}(T-t) \sum_{j=1}^m f_j \sin \lambda_{kn}(t_j-t) \sum_{j=1}^m e_j \cos \lambda_{kn}(t_j-t) \right) \times \\ \quad \times S_{kn}^{-1} \eta_{kn} + v_{kn}(t), \quad 0 \leq t \leq t_1 \\ \left(\sin \lambda_{kn}(T-t) \cos \lambda_{kn}(T-t) \sum_{j=2}^m f_j \sin \lambda_{kn}(t_j-t) \sum_{j=2}^m e_j \cos \lambda_{kn}(t_j-t) \right) \times \\ \quad \times S_{kn}^{-1} \eta_{kn} + v_{kn}(t), \quad t_1 < t \leq t_2 \\ \dots \\ \left(\sin \lambda_{kn}(T-t) \cos \lambda_{kn}(T-t) f_m \sin \lambda_{kn}(t_m-t) e_m \cos \lambda_{kn}(t_m-t) \right) \times \\ \quad \times S_{kn}^{-1} \eta_{kn} + v_{kn}(t), \quad t_{m-1} < t \leq t_m \\ \left(\sin \lambda_{kn}(T-t) \cos \lambda_{kn}(T-t) 0 \ 0 \right) \times \\ \quad \times S_{kn}^{-1} \eta_{kn} + v_{kn}(t), \quad t_m < t \leq t_{m+1} = T \end{cases}$$

Substituting the obtained expressions $u_{kn}(t)$ into (13), we obtain $Q_{kn}(t)$ on the time interval $t \in [0, T]$, and from formulas (7) and (8) we obtain the functions of the deflection, $Q(x, y, t)$, and control, $u(x, y, t)$. Thus, the explicit expressions for the control function $u(x, y, t)$ have the form:

for $0 \leq t \leq t_1$

$$u(x, y, t) = \sum_{k,n=1}^{\infty} \left[\left(\sin \lambda_{kn}(T-t) \cos \lambda_{kn}(T-t) \sum_{j=1}^m f_j \sin \lambda_{kn}(t_j-t) \sum_{j=1}^m e_j \cos \lambda_{kn}(t_j-t) \right) \times \right. \\ \left. \times S_{kn}^{-1} \eta_{kn} + v_{kn}(t) \right] \sin \frac{k\pi}{b} x \sin \frac{n\pi}{c} y,$$

for $t_1 < t \leq t_2$

$$u(x, y, t) = \sum_{k,n=1}^{\infty} \left[\left(\sin \lambda_{kn}(T-t) \cos \lambda_{kn}(T-t) \sum_{j=2}^m f_j \sin \lambda_{kn}(t_j-t) \sum_{j=2}^m e_j \cos \lambda_{kn}(t_j-t) \right) \times \right. \\ \left. \times S_{kn}^{-1} \eta_{kn} + v_{kn}(t) \right] \sin \frac{k\pi}{b} x \sin \frac{n\pi}{c} y,$$

...

for $t_{m-1} < t \leq t_m$

$$u(x, y, t) =$$

$$= \sum_{k,n=1}^{\infty} \left[(\sin \lambda_{kn}(T-t) \quad \cos \lambda_{kn}(T-t) \quad f_m \sin \lambda_{kn}(t_m-t) \quad e_m \cos \lambda_{kn}(t_m-t)) \times \right. \\ \left. \times S_{kn}^{-1} \eta_{kn} + v_{kn}(t) \right] \sin \frac{k\pi}{b} x \sin \frac{n\pi}{c} y,$$

for $t_m < t \leq t_{m+1} = T$

$u(x, y, t) =$

$$= \sum_{k,n=1}^{\infty} \left[(\sin \lambda_{kn}(T-t) \quad \cos \lambda_{kn}(T-t) \quad 0 \quad 0) S_{kn}^{-1} \eta_{kn} + v_{kn}(t) \right] \sin \frac{k\pi}{b} x \sin \frac{n\pi}{c} y.$$

Thus, having explicit expressions for the control function (which is a piecewise continuous function), using the above formulas, the deflection function of the membrane can be found.

CONCLUSION

In this article the problem of control of the vibrations of a rectangular membrane with given non-separated values of the deflection function and velocities at intermediate moments of time by the method of separation of variables, is reduced to the problem of control of countable ordinary differential equations with given initial, final, and non-separated multipoint intermediate conditions. The solution to the problem is constructed by using the methods of the theory of control of finite-dimensional systems with multipoint intermediate conditions.

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ЗАВДАННЯ КЕРУВАННЯ КОЛИВАННЯМИ МЕМБРАНИ З НЕРОЗДІЛЕНИМИ БАГАТОТОЧКОВИМИ УМОВАМИ У ПРОМІЖНІ МОМЕНТИ ЧАСУ

Вступ. Характерною особливістю багатоточкових крайових задач контролю є наявність нерозділених умов у кількох проміжних точках інтервалу дослідження. Такі проблеми керування мають важливе прикладне і теоретичне значення, природно виникає необхідність їхнього дослідження в різних умовах. У статті розглянуто проблему контролю коливань прямокутної мембрани із заданими початковими, кінцевими умовами та нерозділеними значеннями функції перегину та швидкостей у проміжні моменти часу.

Метою статті є розроблення конструктивного підходу до побудови функції керувальної дії для контролю коливань прямокутної мембрани із заданими початковими, кінцевими умовами та нерозділеними (не локальними) значеннями перегину та швидкостями мембранних точок в проміжні моменти часу.

Результати. Методом поділу змінних задача зводиться до задачі керування звичайними диференціальними рівняннями із заданими початковими, кінцевими та нерозділеними багатоточковими проміжними умовами. За допомогою методів теорії керування кінцевомірними системами з багатоточковими проміжними умовами побудовано керувальну дію для контролю коливань прямокутної мембрани.

Висновки. Задача керування коливаннями прямокутної мембрани із заданими нерозділеними значеннями функції відхилення та швидкостей у проміжні моменти часу вирішується за допомогою методів теорії керування кінцево-розмірними системами з багатоточковими проміжними умовами.

Ключові слова: керування коливаннями, коливання мембрани, проміжні значення, нерозділені багатоточкові умови

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**ЗАДАЧА УПРАВЛЕНИЯ КОЛЕБАНИЯМИ МЕМБРАНЫ
С НЕРАЗДЕЛЕННЫМИ МНОГОТОЧЕЧНЫМИ УСЛОВИЯМИ
В ПРОМЕЖУТОЧНЫЕ МОМЕНТЫ ВРЕМЕНИ**

Рассмотрена задача управления колебаниями прямоугольной мембраны с заданными неразделенными значениями функции прогиба и скоростей в промежуточные моменты времени. Методом разделения переменных проблема сводится к задаче управления обыкновенных дифференциальных уравнений с заданными начальными, конечными и неразделенными многоточечными промежуточными условиями. Используя методы теории управления конечномерными системами с многоточечными промежуточными условиями, построено управляющее воздействие.

Ключевые слова: управление колебаниями, колебание мембраны, промежуточные значения, неразделенные многоточечные условия.