

Andreev experiments on superconductor/ferromagnet point contacts

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Andreev reflection is a smart tool to investigate the spin polarization P of the current through point contacts between a superconductor and a ferromagnet. We compare different models to extract P from experimental data and investigate the dependence of P on different contact parameters.

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Following the pioneering work of Igor Yanson and his group at the Institute for Low Temperature Physics and Engineering of the Ukrainian Academy of Sciences (ILTPE NASU) on tiny metallic contacts between two metal electrodes [1], point-contact spectroscopy (PCS) has become a powerful method to study the interactions of ballistic electrons with other excitations in metals [2]. The interpretation of the observed characteristics in point-contact (PC) spectra is usually difficult. These difficulties are frequently inherent in the fabrication of point contacts. In many cases contacts are made by the needle-anvil or shear technique in which two sharpened metal pieces are brought into a gentle touch until a conductive contact is formed. Those contacts are microscopically not well-defined with respect to contact size, geometry, and structure of the metallic nanobridge, and with respect to the local electronic parameters such as the mean free path in the immediate contact region. The only control parameter is the contact resistance, and hence, it is challenging to identify the relevant transport regime free of doubt. Usually Sharvin's [3] or Wexler's [4] formulae for the ballistic and

diffusive transport regime, respectively, are used to infer a PC size estimate from the measured PC resistance. Only recently [5], direct scanning electron microscopy (SEM) measurements of the nanocontact size of nanostructured point contacts allowed for the first time a direct comparison with theoretical models for contact-size estimates of heterocontacts. The semiclassical models yield reasonable values for the PC radius a as long as the correct transport regime is determined by taking into account the local transport parameters of the individual contact. Of course, this requires a careful characterization of the samples with respect to the local resistivity and the local mean free path.

Among the rich variety of solid-state problems investigated by point-contact spectroscopy the study of superconductor-metal contacts contributes a significant portion. Nowadays point-contact spectroscopy is an important tool to explore the symmetry and nodal structure of the energy gap Δ of conventional and unconventional superconductors [6]. When the temperature is lowered below the superconducting transition temperature T_c of the superconducting electrode of a superconductor (S) / normal metal (N)

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point-contact Andreev reflection [7] of charge carriers at the S/N interface occurs. Andreev reflection leads to minima at $V \approx \pm\Delta/e$ in the differential resistance dV/dI as a function of applied bias V , i.e., maxima in the corresponding conductance curves $G(V) = (dI/dV)(V)$, and thus allows determination of the gap size, also while varying temperature and magnetic field, respectively [8,9].

A new pitch came into the field when Andreev reflection was used to extract the spin polarization P of the current through superconductor/ferromagnet (F) point contacts [10,11]. Knowledge of the spin polarization of possible materials for spin-electronic devices is a key issue for spintronics [12]. An efficient spin injection is of central importance for utilizing the spin degree of freedom as a new functionality in spin-electronic devices. The spin polarization P of ferromagnets can be measured by various techniques including photoemission [13], spin-dependent tunnelling [14], and point-contact Andreev reflection (PCAR) [15]. For a quantitative analysis of the results, however, one has to be aware of the different nature of the quantities measured by each technique. The spin polarization, defined by the difference of spin-up and spin-down density-of-states, is typically measured with spin-polarized photoemission while the spin polarization of the transport current is obtained, e.g., in PCAR experiments [16], which in turn is distinctly different from the spin polarization of the density-of-states resulting from tunneling experiments [14]. An issue of considerable importance is how the spin polarization obtained by Andreev reflection is related to the ferromagnet's bulk spin polarization [17].

Already a variety of materials have been investigated including the ferromagnetic elements Fe, Co, and Ni and several alloys mainly with Al, Nb, or Pb as a superconducting counter electrode [10,11,18–23]. However, different models [10,11,24–26] describing the transport through S/F interfaces yielded varying values for P , also depending on the contact fabrication and the transport regime [27], an issue that is not yet understood in detail [28]. In the following, we want to review the main ideas of two most prominent models shortly and compare the results of both to the same set of experimental data obtained on nanostructured Al/Fe contacts [23].

The theoretical analysis of most S/F point-contact experiments has been carried out in the spirit of the Blonder–Tinkham–Klapwijk (BTK) theory [9] for Andreev reflection at an interface between N and classical S with spin-singlet pairing. This is the coherent process by which an electron from N enters S and a hole of opposite spin is retro-reflected, creating a spin-singlet Cooper pair in S. Possible ordinary reflection at the S/F interface barrier is parametrized by a phenomenological parameter, the barrier strength Z . The sensitivity of the Andreev process to the spin of the carriers originates from the conservation of the spin direction at the interface. Consequently, when there is an imbalance in the number of spin-up and spin-down elec-

trons at the Fermi level, as it is the case in the spin-polarized situation of a ferromagnetic metal, this leads to a reduction of the Andreev reflection probability [15]. Andreev reflection is limited by the minority carriers of the metal.

In the simplest approach applied for the analysis of several experiments [10,18–21,24], the total current I through the constriction is decomposed into a fully unpolarized part $(1-P')I_u$ for which Andreev reflection is allowed and into a fully polarized part $P'I_p$ for which Andreev reflection is zero,

$$I = (1-P')I_u + P'I_p.$$

The weighting factor P' determines the spin polarization of the ferromagnet. In the following we will refer to this model as the dispartment model. Consequentially, the conductance G_{SF} is also decomposed into two parts:

$$G_{SF} = (1-P')G_u(V) + P'G_p(V)$$

with

$$G_{u,p}(V) = \int_{-\infty}^{\infty} \frac{df(E-V, T)}{dV} [1 + A_{u,p}(E, Z) - B_{u,p}(E, Z)] dE,$$

where G_u denotes the conductance, A_u the Andreev reflection probability, and B_u the normal reflection probability of the fully unpolarized channel, and G_p , A_p , and B_p denote the corresponding quantities of the fully polarized channel. Both contributions are derived in the BTK formalism and following expressions for the zero-temperature conductances G_u and G_p are obtained [25]:

	$ E < \Delta$	$ E > \Delta$
$G_u _{T=0}(E)$	$\frac{2(1+\beta^2)}{\beta^2 + (1+2Z^2)^2}$	$\frac{2\beta}{1+\beta+2Z^2}$
$G_p _{T=0}(E)$	0	$\frac{4\beta}{(1+\beta)^2 + 4Z^2}$

with $\beta = E / \sqrt{|\Delta^2 - E^2|}$.

Despite the attractive simplicity of the BTK formalism it has been shown [17,27] that application of the BTK formalism (even in its generalized form [25]) has certain drawbacks and enforces several assumptions for the analysis. This has mainly to do with the problem to determine P' and Z independently. The physical reason is that both lead to a reduction of the Andreev current and diminish the conductance change in $G = dI/dV$. The model fails to distinguish whether it is high P or high Z that causes the depression of conductance at small bias. This problem is evaded by applying a different theoretical approach [26].

Cuevas and coworkers [22,26] developed a model based on quasiclassical Green functions. The central quantities of the model are two transmission coefficients $\tau_{\uparrow,\downarrow} = |t_{\uparrow,\downarrow}|^2$. Therefore, we will refer to this model as the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model throughout this paper. The transmission coefficients contain all microscopic properties relevant for the transport through the constriction, i.e., they account for the majority- and minority-spin bands in the ferromagnet, the electronic structure of the superconductor, and the interface. $t_{\uparrow,\downarrow}$ and $r_{\uparrow,\downarrow} = (1 - \tau_{\uparrow,\downarrow})^{1/2}$ are the spin-dependent transmission and reflection amplitudes, respectively, entering the normal-state scattering matrix \hat{S} which supplements the boundary conditions of the theory. Of course, the restriction to a single conduction channel per spin direction is a rough simplification of the point contact, but it is finally justified by the agreement with the experiment [22,23]. Following the calculation by Cuevas and coworkers the spin-dependent current through the S/F point contact can be separated in two spin contributions,

$$I_{SF} = I_{\uparrow} + I_{\downarrow}$$

and each contribution can be written in its BTK form [9]

$$I_{\sigma} = \frac{e}{h} \int_{-\infty}^{\infty} d\varepsilon [f(\varepsilon - eV) - f(\varepsilon)] [1 + A_{\sigma}(\varepsilon) - B_{\sigma}(\varepsilon)],$$

where $f(E)$ is the Fermi function, $A_{\sigma}(\varepsilon)$ and $B_{\sigma}(\varepsilon)$ are the spin-dependent Andreev reflection and normal reflection probabilities, respectively, and $\sigma = \uparrow$ or \downarrow . $A_{\sigma}(\varepsilon)$ ($B_{\sigma}(\varepsilon)$) is calculated from the spin-dependent transmission (reflection) amplitudes, and finally, the zero-temperature conductance of the S/F contact adopts the form

$$G_{SF} = \frac{4e^2}{h} \begin{cases} \frac{\tau_{\uparrow}\tau_{\downarrow}}{(1+r_{\uparrow}r_{\downarrow})^2 - 4r_{\uparrow}r_{\downarrow}(eV/\Delta)^2}, & eV \leq \Delta \\ \frac{\tau_{\uparrow}\tau_{\downarrow} + (\tau_{\uparrow} + \tau_{\downarrow} - \tau_{\uparrow}\tau_{\downarrow})\sqrt{1 - (\Delta/eV)^2}}{[(1-r_{\uparrow}r_{\downarrow}) + (1+r_{\uparrow}r_{\downarrow})\sqrt{1 - (\Delta/eV)^2}]^2}, & eV > \Delta \end{cases}$$

while the normal-state conductance is given by

$$G_N = \frac{e^2}{h} (\tau_{\uparrow} + \tau_{\downarrow}).$$

It is obvious that the Andreev spectra are determined by a set of three free parameters τ_{\uparrow} , τ_{\downarrow} , and Δ . The current spin polarization P in this model is defined by

$$P = \frac{|\tau_{\uparrow} - \tau_{\downarrow}|}{\tau_{\uparrow} + \tau_{\downarrow}}$$

and can be determined from the fit parameters of an experimental Andreev spectrum. We note, that this expression is symmetric with respect to τ_{\uparrow} and τ_{\downarrow} , therefore, one cannot assign a transmission coefficient to the majority or mi-

nority charge carriers in the ferromagnet. However, we expect the high transmissive coefficient τ_{\downarrow} to correspond to the minority electrons. In the absence of spin polarization, i.e., $P=0$ for a N/S contact with $\tau_{\uparrow} = \tau_{\downarrow}$, above formulae reduce to the well-known BTK result [9].

Figure 1 displays a set of normalized conductance curves for $T = 0.02T_c$, $\tau_{\downarrow} = 0.99$, and $\tau_{\uparrow} = 0.99, 0.53, 0.42, 0.33, 0.25$, and 0.17 which equals $P = 0, 0.3, 0.4, 0.5, 0.6, 0.7$ from top to bottom. The shape of each spectrum is unambiguously determined by a set of τ_{\uparrow} , τ_{\downarrow} . For $P = 0$ (top curve) the curve reproduces the well-known BTK result where Andreev reflection causes a doubling of the normal-state conductance for energies $eV < \Delta$. We note that the characteristic double-peak feature at $|eV| = \Delta$ caused by the reduction of the conductance at low bias originates from a small fraction of charge carriers ordinarily reflected at an interface barrier. In the case of two spin-dependent transport channels it is intuitive to consider the high-transmissive spin channel to get a measure for the fraction of charge carriers ordinarily reflected at the interface barrier. From the transmission coefficients $\tau_{\downarrow} = 1/(1 + Z_{\downarrow}^2)$ one derives $Z_{\downarrow} = ((1 - \tau_{\downarrow})/\tau_{\downarrow})^{1/2}$. For the upper curve which corresponds to the unpolarized BTK case one gets $Z_{\downarrow} = 0.1$ as the parameter equivalent to the BTK interface parameter Z .

While the shape of the spectra is unambiguously determined by a set of τ_{\uparrow} , τ_{\downarrow} , the spectra for same polarization can look quite different. Figure 2 shows four curves calculated for different sets of τ_{\uparrow} , τ_{\downarrow} which all result in $P = 0.4$. The calculations have been performed for a finite temperature $T = 0.22T_c$. The high-transmission spin-channel seems to be decisive whether the curve shape appears more point-contact-like or more tunnelling-like.

Before we compare curves calculated by both models, let us first check the validity of the fitting procedure. For this purpose we measured point-contact spectra of S/N point contacts in a ^4He cryostat down to $T = 1.4$ K and

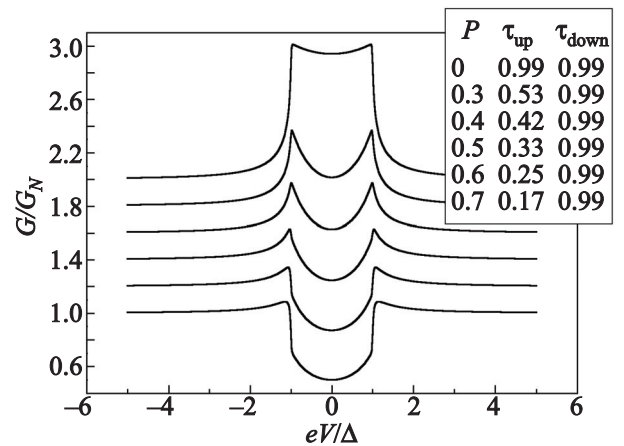


Fig. 1. Normalized differential conductance G/G_N vs. eV/Δ . The curves are calculated with the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model (see text) for different polarizations at $T = 0.02T_c$. For clarity, the curves are shifted successively upwards by 0.2 units.

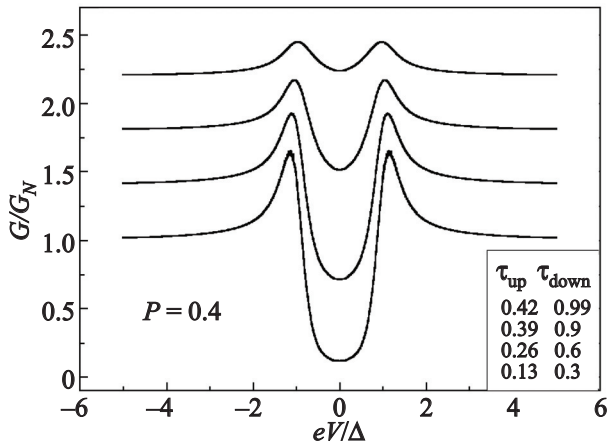


Fig. 2. Normalized differential conductance G/G_N vs. eV/Δ . The curves are calculated with the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model (see text) for the same polarization $P=0.4$ at $T=0.22T_C$. For clarity, the curves are shifted successively upwards by 0.4 units.

fitted them with both models. The PCs have been established in the edge-to-edge configuration with a sharpened Pb electrode and a normal-metal electrode made from Cu or Pt. Therefore, for both fits we expect $P' = P = 0$. Fig-

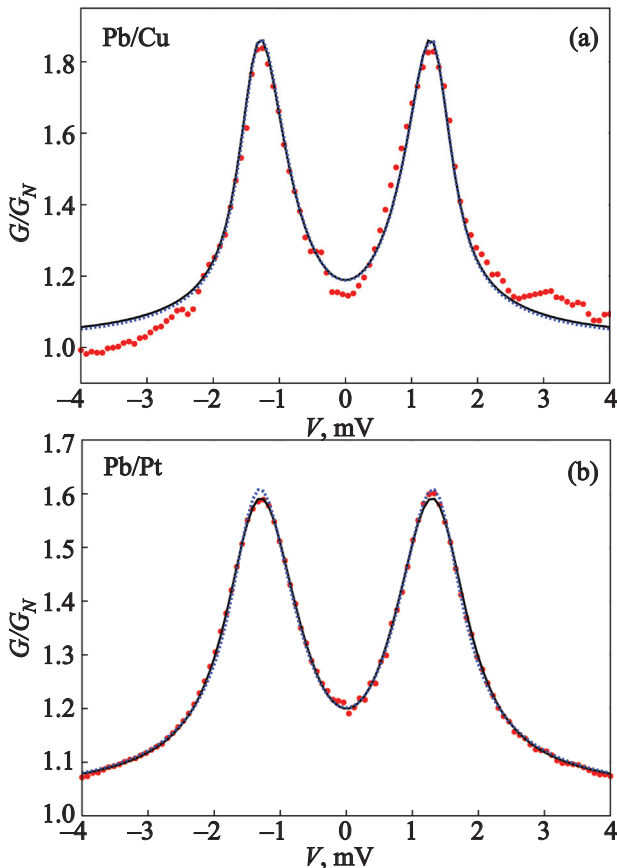


Fig. 3. (Color online) Point-contact spectra of S/N contacts together with fits according to Martin–Rodero *et al.* [26] (dashed line) and Mazin *et al.* [25] (solid line). (a) Pb/Cu contact at 1.4 K with $R = 2.7 \Omega$. (b) Pb/Pt contact at 2.1 K with $R = 7.5 \Omega$. For the fit parameters see text.

ures 3(a) and (b) show the experimental data together with the fits according to the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model [26] (dashed line) and the dispartment model [25] (solid line). In both cases we got almost perfect agreement with the data. For the Pb/Cu contact with $R = 2.7 \Omega$ measured at $T = 1.4$ K (upper panel) we obtained $\Delta_0 = \Delta(T = 0) = 1.38$ meV, $\tau_{\uparrow} = \tau_{\downarrow} = 0.814$ as parameters for the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model, and $\Delta_0 = 1.38$ meV, $Z = 0.475$, $P' = 0.011$ for the dispartment model. For the Pb/Pt contact with $R = 7.5 \Omega$ measured at $T = 2.1$ K (lower panel) the parameters are $\Delta_0 = 1.48$ meV, $\Gamma/\Delta_0 = 0.10$, $\tau_{\uparrow} = 0.811$, $\tau_{\downarrow} = 0.819$ for the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model, and $\Delta_0 = 1.42$ meV, $\Gamma/\Delta_0 = 0.10$, $Z = 0.488$, $P' = 0.015$ for the dispartment model. In both cases we found within the experimental error a good agreement of $P \approx P' \approx 0$ as expected for these non-magnetic metals. The energy gap of Pb determined from the fits coincides fairly good to the gap value reported in literature [29]. We note that a small broadening of the Pb/Pt spectra caused by inelastic scattering in the contact region is accounted for by introducing the Dynes [30] parameter Γ which is of the order of 5–10% of Δ_0 .

In the next step we compare curves calculated with the simple dispartment model to those calculated with the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model for nominal same polarization $P = P'$. Figure 4 displays a set of normalized conductance curves calculated for $T = 0.1$ K, $P = P' = 0, 0.2, 0.4, 0.6, 0.8$, and 1, and $Z = 0.3$ for the dispartment model, and $\tau_{\downarrow} = 0.917$ for the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model, respectively, which corresponds to $Z_{\downarrow} = [(1 - \tau_{\downarrow})/\tau_{\downarrow}]^{1/2} = 0.3$. In the extreme case $P = P' = 0$ which describes a S/N contact the calculations perfectly agree with each other, as well as for the other extreme case $P = P' = 1$ which describes a halfmetallic S/F contact. For the latter, there is only a small difference in the vicinity of the coherence peaks at $|eV/\Delta| = 1$. However, at intermediate values with increasing polarization the Andreev signal is much faster suppressed for the dispartment model

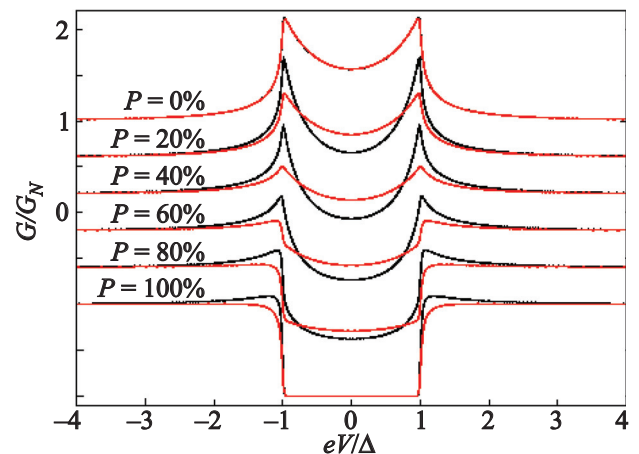


Fig. 4. (Color online) Comparison of Andreev spectra with $P = 0, 0.2, 0.4, 0.6, 0.8$, and 1 calculated according to the $\tau_{\uparrow}\text{-}\tau_{\downarrow}$ -model (black lines) and the dispartment model (red lines). For clarity, the curves are shifted successively downwards by 0.4 units.

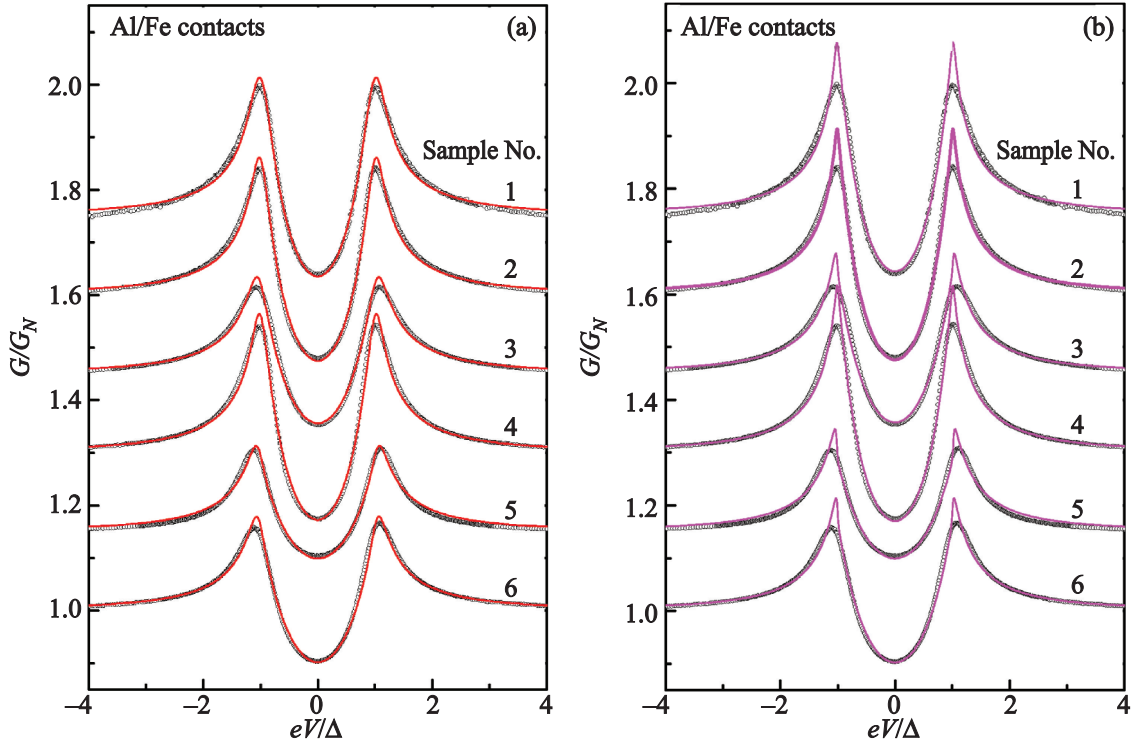


Fig. 5. (Color online) Point-contact spectra of Al/Fe contacts at 0.1 K together with fits according to the τ_{\uparrow} - τ_{\downarrow} -model (a) and according to the dispartment model (b).

than for the τ_{\uparrow} - τ_{\downarrow} -model. Our comparison discloses a notable difference of both quantities which makes questionable a contrasting juxtaposition of P and P' values derived from the analysis of experimental data by one of these models.

In order to illustrate this difference on experimental data we used both models to fit the same set of data measured on Al/Fe nanostructured point contacts [23]. The PCs were fabricated by structuring a hole of 5–10 nm diameter with electron-beam lithography and subsequent reactive ion etching into a 50-nm thick Si_3N_4 membrane, evaporating on one side a 200 nm Al layer, and on the other side a 12 nm thick Fe layer and a Cu layer of thickness $d_{\text{Cu}} = 188$ nm as a low-ohmic electrode [31]. Figure 5 displays the normalized conductance spectra of six different nanostructured PCs with contact resistances between 2.7 and 24.2 Ω measured in a dilution refrigerator at

$T = 0.1$ K together with fitting curves calculated with the τ_{\uparrow} - τ_{\downarrow} -model (a) and the dispartment model (b). Both models perfectly describe the data, minor deviations are observed only at $|eV/\Delta| \approx 1$ where the experimental curves are more rounded probably caused by a leveling-off of the electron temperature due to heating by electromagnetic stray fields or a small pair-breaking effect by Fe. The corresponding fit parameters are listed in Table 1. Within an uncertainty of 1% the same gap value Δ_0 is found for both models, however, there is a notable difference in the Z parameter, which is a factor 2–3 higher in the dispartment model, and the spin polarization P' which is lower. Although the origin of Z is not clear at all, in both models it subsumes all ordinary reflection of charge carriers that occurs at the interface for $P = P' = 0$, e.g., reflection caused by an insulating interface layer, lattice imperfections, Fermi velocity mismatch, etc. For P and $P' > 0$

Table 1. Fit parameters Δ_0 , τ_{\uparrow} , and τ_{\downarrow} of the τ_{\uparrow} - τ_{\downarrow} -model and Δ_0 , Z , and P' of the dispartment model for conductance spectra of nanostructured Al/Fe point contacts. From the transmission coefficients τ_{\uparrow} and τ_{\downarrow} the interface barrier Z_{\downarrow} and the current spin polarization P have been calculated.

Sample No.	R_N, Ω	Δ_0, meV	τ_{\uparrow}	τ_{\downarrow}	Z_{\downarrow}	P	Δ_0, meV	Z	P'
1	2.68	0.174	0.371	0.983	0.132	0.452	0.176	0.337	0.396
2	6.98	0.175	0.362	0.984	0.128	0.460	0.176	0.338	0.407
3	7.29	0.157	0.349	0.993	0.083	0.480	0.158	0.272	0.444
4	9.59	0.190	0.361	0.984	0.128	0.463	0.191	0.342	0.407
5	18.4	0.166	0.348	0.997	0.054	0.482	0.168	0.218	0.497
6	24.2	0.174	0.343	0.994	0.078	0.487	0.174	0.262	0.494

the situation is less apparent. The physical reason is that Z and P' both lead to a reduction of the Andreev current. The dispartment model obviously fails to distinguish whether a high-spin polarization or a high barrier causes the depression whereas for the $\tau_{\uparrow}-\tau_{\downarrow}$ -model as per definition only the conductance channel not affected by the suppression of Andreev reflection due to polarization is considered to determine ordinary reflection.

Another important aspect is that the previously reported dependence of the spin polarization on the contact size [23] is robust against the model used to extract P . Independently of the model there is a clear reduction of the spin polarization with decreasing contact resistance R_N , i.e., increasing contact radius a . The reduction of P has been allocated as being due to spin-orbit scattering in the contact region with a constant scattering length ℓ_{so} modelled by a simple exponential decay [23], $P(a) = P_0 \exp(-a/\ell_{so})$. A spin-orbit scattering length $\ell_{so} = 255$ nm has been obtained for the analysis with the $\tau_{\uparrow}-\tau_{\downarrow}$ -model. The same systematic trend of $P'(a)$ is found for the dispartment model albeit with a lower value for ℓ_{so} . For small contacts both models result in almost the same spin polarization.

In summary, we have discussed possible reasons for the scatter of polarization values found for the spin polarization measured by Andreev reflection in point-contact experiments. We showed that the scatter is partially caused by the models used to extract the spin polarization from the data, and partially caused by intrinsic mechanisms like the spin-orbit scattering in the contact region.

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